

1

Work in preparation:

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on RELATIVE PARTITION FUNCTION
FOR COULOMB + POINT INTER.

$$H_0 = -\Delta + \frac{\gamma}{|x|}$$

in $L^2(\mathbb{R}^3)$, $\gamma \in \mathbb{R}$

Well def., s.a., lower bounded

$$R(\lambda; H_0)(x, y) = - \frac{\Gamma\left(1 - \frac{\gamma}{2\sqrt{-\lambda}}\right)}{4\pi |x-y|}$$

$K\left(-\frac{\gamma}{2\sqrt{-\lambda}}; \sqrt{-\lambda} x_+, \sqrt{-\lambda} x_-\right)$, with

$$K(a; b, c) := W_{a, \frac{1}{2}}(b) M'_{a, \frac{1}{2}}(c) - \\ - W'_{a, \frac{1}{2}}(b) M_{a, \frac{1}{2}}(c)$$

$\lambda \in \rho(H_0)$, $\operatorname{Re} \sqrt{-\lambda} > 0$, $x, y \in \mathbb{R}^3$

$$x_{\pm} := |x| + |y| \pm |x-y|$$

M, W : Whittaker functions

$$\Gamma(s) := \int_0^{+\infty} x^{s-1} e^{-x} dx, \operatorname{Re} s > 0 \dots$$

Prop. 1 i) $\sigma_{\text{ess}}(H_0) = \sigma_{\text{ac}}(H_0) = [0, +\infty)$

ii) If $\gamma \geq 0$: $\sigma_{\text{pp}}(H_0) = \emptyset$

If $\gamma < 0$: $\sigma_{\text{pp}}(H_0) = -\frac{\gamma^2}{4(n+1)^2}$,
 $n \in \mathbb{N}_0$

$$H_\alpha = \text{"} -\Delta + \frac{\gamma}{|\cdot|} + \alpha \delta_0 \text{"}$$

$-\infty < \alpha \leq +\infty, \gamma \in \mathbb{R}$.

rigorously given as lower bounded, s.a. operator by "resolvent kernel"

$$R(\lambda; H_\alpha)(x, y) = R(\lambda; H_0)(x, y)$$

$$= \frac{4\pi}{4\pi\alpha - \gamma F_\gamma\left(\frac{\gamma}{2\sqrt{-\lambda}}\right)} g(\lambda; x) g(\lambda; y),$$

$x, y \in \mathbb{R}^3$

$$\lambda \in \rho(H_\alpha) \cap \rho(H_0), \operatorname{Re} \sqrt{-\lambda} > 0$$

$$g(\lambda; x) := \frac{\Gamma\left(1 + \frac{\gamma}{2\sqrt{-\lambda}}\right)}{4\pi|x|} W_{\left(2\sqrt{-\lambda}|x|, -\frac{\gamma}{2\sqrt{-\lambda}}, \frac{1}{2}\right)} \quad x \neq 0$$

$$F_{\gamma}(z) := \begin{cases} \Psi(1+z) - \log z - \frac{1}{2z} - \Psi(1) - \Psi(2), & \gamma > 0 \\ \Psi(1+z) - \log(-z) - \frac{1}{2z} - \Psi(1) - \Psi(2), & \gamma < 0, \end{cases} \quad z \in \mathbb{C}$$

$$\Psi(z) := \frac{d}{dz} \log \Gamma(z), \quad z \in \mathbb{C}$$

"digamma function"

Prop. 2 i) $\sigma_{\text{ess}}(H_{\alpha}) = \sigma_{\text{ac}}(H_{\alpha}) = [0, +\infty)$

ii) Eigenvalues $E < 0$ of H_{α}
"for $l=0$ " solve

$$4\pi\alpha = \gamma F_{\gamma}\left(\frac{\gamma}{2\sqrt{-E}}\right)$$

iii) For $\gamma \geq 0$, $\alpha > -\gamma [\psi(1) + \psi(2)]/4\pi$

$$\sigma_{pp}(H_\alpha) = \emptyset$$

iv) For $\gamma \geq 0$, $\alpha < -\gamma [\psi(1) + \psi(2)]/4\pi$

$$|\sigma_{pp}(H_\alpha)| = 1$$

(eigenvalue given by above eqt.)

v) $\gamma < 0$: ∞ -many solutions of above eqt., simple eigenvalues corresp. to $l=0$

For $l \geq 1$: Eigenvalues
 $-\frac{\gamma^2}{4m^2}$, $m \geq 2, m \in \mathbb{N}$

Consider now the case $\gamma \geq 0$

Prop. 3 The "relative trace"

$r(\lambda; H_\alpha, H_0)$ is well def. for
 $\lambda \in \rho(H_0) \cap \rho(H_\alpha)$,

$$r(\lambda; H_\alpha, H_0) = \text{Tr} [R(\lambda; H_\alpha) - R(\lambda; H_0)]$$

One has:

$$r(\lambda; H_\alpha, H_0) = - \frac{z [2\psi'(1+z)z^2 - 2z + 1]}{A(\gamma, \alpha, z)},$$

$$\text{with } A(\gamma, \alpha, z) := \gamma(4\pi\alpha - \gamma)(\psi(1+z) - \log z - \frac{1}{2z} - \psi(1) - \psi(2)),$$

$$\text{with } \operatorname{Re} \sqrt{-\lambda} > 0, \quad z := \frac{\gamma}{2\sqrt{-\lambda}}.$$

Proof: Use resolvent formula, integration, Γ -functional equation...

Rem: Using asympt. exp. of digamma-function ψ one gets:

$$r(\lambda; H_\alpha, H_0) = \sum_{k=0}^{\infty} b_k (-\lambda)^k, \quad b_k \in \mathbb{R}, \quad |\lambda| \ll 1$$

$$e^{-\gamma} b_0 = [3\gamma(\gamma - 2C\gamma + 4\pi\alpha)]^{-1}$$

$$b_1 = \frac{(17 - 24C)\gamma + 48\pi\alpha}{45\gamma^3(2C\gamma - \gamma - 4\pi\alpha)^2}$$

$$\sum_{j=2, k=0}^{\infty} a_{j,k} (-\lambda)^{-\frac{1}{2}} (\log(-\lambda))^k, \quad |\lambda| \gg 1$$

$$a_{j,k} \in \mathbb{R}$$

$$a_{2,0} = -\frac{1}{2}, \quad a_{2,k>0} = 0,$$

$$a_{3,0} = \frac{1}{2} [4\pi\alpha + (2-C)\delta + \delta (\log\delta - \log 2)]$$

$$a_{3,1} = -\frac{\delta}{4}$$

$$a_{3,k>1} = 0.$$

One gets corresponding formulae for "relative spectral measure

$$e(\nu; H_\alpha, H_0), \quad \nu \in (0, +\infty),$$

s.t. , e.g.

$$\text{Tr} (e^{-tH_\alpha} - e^{-tH_0}) =$$

$$= \int_0^\infty e^{-\nu^2 t} e(\nu; H_\alpha, H_0) d\nu$$

Namely:

$$e(\nu; H_\alpha, H_0) = \frac{\nu}{i\pi} \lim_{\varepsilon \downarrow 0+} \left[\pi (v^2 e^{2\pi i - i\varepsilon}; H_\alpha, H_0) - \pi (v^2 e^{i\varepsilon}; H_\alpha, H_0) \right],$$

with "explicit expansions" for $\nu \downarrow 0$, $\nu \uparrow +\infty$.